Magnetic Microstructure Aluminum-substituted Barium Hexaferrite for Microwave Devices mm-Wavelength Range

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The magnetic structure and cationic distribution was investigate in ferrite with a hexagonal structure $BaFe_{12-x}Al_xO_{12}$ ($x \le 2$). It is found that when x = 2 Al³⁺ ions are localized mainly in the positions of 12k (*a*-sublattice) and 4f₁ (*c*-sublattice). This contribute to change of number magnetically active connections and therefore the magnetic structure of ferrite and magnetic parameters. It was concluded about the possibility of using the results obtained with the choice of optimal composition of substrates for millimeterwave devices.

Keywords: Hexaferrite, Magnetic microstructure, Cationic distribution, Microwave, Crystal structure.

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1. INTRODUCTION

Currently the problem of nano architectonic come to the forefront at obtaining new crystalline materials for microwave magnetoelectronic. They include the production of thin layers of magnetic materials functioning in weak magnetic fields, or lack thereof. Such layers can be obtained by ordering of nanoscale structural elements as a result of coordinated interaction between the ferromagnetic nanostructures and solid surface (substrate), has a strong internal fields of the magnetic anisotropy. Compliance with this requirement is only possible if a certain composition, structure and physical properties of the surface of the magnetic substrate.

A very promising material for such substrates are ferrites with a hexagonal structure [1]. They have a valuable combination of magnetic and dielectric properties capable of providing the required level of parameters required for the creation of microwave devices [2]. First of all, for their peculiar presence of high values of the magnetic anisotropy, adjustable level of saturation magnetization and dielectric loss. The crystal structure executes the important role in the formation of these propererties. The crystal structure includes hexagonal layers, in which a portion of the oxygen ions is substituted with ions Ba²⁺, Sr²⁺, Ca²⁺, etc. [3]. These layers provide a preferential orientation of the atoms of iron ions, and iron ions located in five-fold oxygen coordination, build a strong internal magnetic anisotropy field [4].

However, investigations of the hexagonal ferrites with different types of characterizing the cations (Ba, Sr, Ca) indicate the weak influence of cations on the energy of the magnetic anisotropy and the magnetization of the hexagonal ferrites. A more significant effect on these parameters have diamagnetic atoms replacing Fe ions and change the number and strength of the exchange coupling between the five existing magnetic sublattices. In addition, they are able to change the local distortions of the crystal lattice and consequently, reduce internal stresses and affect the magnetic structure of matter. These ion are In^{3+} , Sc^{3+} , Al^{3+} and others [1, 5].

The paper studied the effect of Al^{3+} ions on the magnetic structure and magnetic properties of hexagonal ferrite $BaFe_{12-x}Al_xO_{19}$.

The chemical composition of of ferrites was set based on the presence of magnetic anomalies in the concentration dependences of the magnetic parameters and data of Mossbauer studies $BaFe_{12-x}Al_xO_{12}$ [3, 5].

2. EXPERIMENTAL PART

The objects for the study were anisotropic and isotropic polycrystalline of ferrites $BaFe_{12}O_{19}$ and $BaFe_{10}Al_2O_{19}$. The ferrites were manufactured by the known ceramic technology and by X-ray analysis, was not contain external phases. The magnetic parameters: specific saturation magnetization σ_8 , specific residual magnetization σ_7 , coercive force H_c , Curie temperature, anisotropy field H_a measured by standard techniques on a vibrating magnetometer VSM250 in the field of 20 kOe. Mössbauer investigations were conducted on the spectrometer Ms 1104 Em with automatic processing the spectra of the program Univem Ms in the range $87\div300$ K.

3. RESULTS AND DISCUSSION

Fig. 1 show as an example Mössbauer spectra of isotropic and anisotropic of ferrites samples $BaFe_{12}O_{19}$ and $BaFe_{10}Al_2O_{19}$.

From Fig. 1a and b shows that the spectra of samples $BaFe_{12}O_{19}$ represent a superposition of five sextets corresponding sublattices a (12 κ), b (2a), c (4f1), d (4f2), e (2B). Population sublattices correspond to their theoretical values for the hexagonal structure of M. The parameters of Mössbauer spectra of iron ions localized in these sublattice coincide with the data obtained in [6].

In contrast to the unsubstituted hexaferrites $BaFe_{12}O_{19}$ Mössbauer spectra of samples of substituted of ferrites $BaFe_{12} - _xAl_xO_{12}$ represent a superposition of 7 sextets, corresponding sublattices a_1 , a_2 , b, c_1 , c_2 , d

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and *e*. The emergence of non-equivalent nodes $12 \text{K m} 4f_1$ associated with the peculiarities of the localization Al³⁺. According to the [3] it is assumed that the parameters of the substitution of Fe³⁺ ions Al³⁺ ions (x < 5) Al³⁺ ions localizes in positions 12 K (*a*-sublattice), 2a (B-sublattice) or $4f_1$ (*c*-sublattice). It can be assumed that the Al³⁺ ions are localized mainly in the positions 12 K and $4f_1$ comparing area sextets (Table. 1).



Fig. 1 – Mössbauer spectra of Ba-ferrites: a – isotropic BaFe₁₂O₁₉; b – anisotropic BaFe₁₂O₁₉; c – isotropic BaFe₁₀ Al₂O₁₉; d – anisotropic BaFe₁₀ Al₂O₁₉

This assumption is confirmed by the decrease in the local magnetic fields (H_L) to sublattices d (4f₂) and c (2_{B1}). The reduction values H_L in these sublattices indi-

cates a weakening of the main exchange relations, due to the substitution of Fe³⁺ ions ions Al³⁺. From an examination of circuit of main exchange interactions follows that the decrease H_{eff} for sublattices d and erelatively unsubstituted ferrite BaFe₁₂O₁₉ associated with preferential localization of Al³⁺ ions in 12 κ and 4f₁positions. Indeed, the sublattice c has the maximum number of (9) relations with a-under the grate. In turn, sublattice a has exchange connection with c-sublattice and 2 connection with d-sublattice.

Loss of magnetically bonds Fe_c -O-Fe_a and Fe_a -O-Fe_d leads to a decrease local magnetic fields at the nucleus Fe^{57} , the appearance of non-equivalent nodes 12k and 4f₁ and as a consequence to a reduction sublattice magnetization and other parameters of the ferrite (Table 2).

From the Table 2 it is seen that the weakening of the exchange coupling leads to a decreases the Curie temperature of 177 °C, the specific saturation magnetization of 43 Am²/kg compared with BaFe₁₂O₁₉ ferrite. However, based on the cationic distribution calculated reduction of the specific saturation magnetization of ferrite is substantially lower than observed experimentally. This indicates the presence of non-collinear magnetic structure in the ferrite BaFe10Al2O19. Based on the ratio of integral intensities of the components of the Mössbauer spectra of anisotropic samples, it was found that the angle of deviation of the magnetic moments of the sublattices of ferrite is 44.5 degrees. The deflection angle of the magnetic moments retains its value, and the temperature is lowered to 87 K. Non-collinear magnetic structure is accompanied by a change in the shape of the magnetic hysteresis loop. As an example, Fig. 2 shows the magnetic hysteresis loop ferrite BaFe12O19 and BaFe10Al2O19.



Fig. 2 – Hysteresis loops anisotropic hexagonal ferrites: a – BaFe₁₂O₁₉; b – BaFe₁₀Al₂O₁₉

| Ferrite | Sublattice, | The iso- | Quadrupole | Magnetic field at | Areas |
|---|---|-------------------|----------------------|------------------------------|----------------|
| | component | mer shift | splitting Δ , | the nucleus Fe ⁵⁷ | component S, % |
| | spectrum | δ , mm/sec | mm/sec | H, kOe | |
| | | | | | |
| | | | | | |
| Isotropic BaFe12O19 powder | <i>a</i> -C1(Fe ³⁺)VI | 0,35 | 0,42 | 414 | 49,9 |
| | <i>b</i> -C4(Fe ³⁺)VI | 0,33 | 0,00 | 507 | 8,8 |
| | <i>c</i> -C3(Fe ³⁺)IV | 0,26 | 0,22 | 491 | 18,6 |
| | <i>d</i> -C2(Fe ³⁺)VI | 0,39 | 0,20 | 516 | 17,6 |
| | e-C5(Fe ³⁺)V | 0,29 | 2,21 | 400 | 5,1 |
| Anisotropic BaFe12O19 powder | <i>a</i> -C1(Fe ³⁺)VI | 0,35 | 0,41 | 414 | 47,7 |
| | <i>b</i> -C4(Fe ³⁺)VI | 0,32 | 0,02 | 506 | 7,8 |
| | <i>c</i> -C3(Fe ³⁺)IV | 0,26 | 0,21 | 488 | 19,4 |
| | <i>d</i> -C2(Fe ³⁺)VI | 0,39 | 0,18 | 514 | 19,4 |
| | e-C5(Fe ³⁺)V | 0,28 | 2,17 | 403 | 5,7 |
| Isotropic BaFe ₁₀ Al ₂ O ₁₉ powder | <i>a</i> ₁ - C1(Fe ³⁺)VI | 0,31 | 0,38 | 370 | 18,0 |
| | <i>d</i> -C2(Fe ³⁺)VI | 0,37 | 0,13 | 474 | 22,9 |
| | c_1 -C3(Fe ³⁺)IV | 0,31 | 0,43 | 397 | 23,2 |
| | <i>b</i> -C4(Fe ³⁺)VI | 0,37 | 0,52 | 311 | 8,2 |
| | a_{2} - C5(Fe ³⁺)IV | 0,28 | 0,24 | 428 | 16,6 |
| | c_{2} - C6(Fe ³⁺)IV | 0,22 | 0,37 | 345 | 2,6 |
| | e-C7(Fe ³⁺)V | 0,29 | 2,1 | 383 | 8,4 |
| Anisotropic BaFe ₁₀ Al ₂ O ₁₉ foil | a_1 - C1(Fe ³⁺)VI | 0,29 | 0,39 | 396 | 25,8 |
| | <i>d</i> -C2(Fe ³⁺)VI | 0,37 | 0,13 | 485 | 17,4 |
| | c_1 -C3(Fe ³⁺)IV | 0,25 | 0,32 | 423 | 14,8 |
| | <i>b</i> -C4(Fe ³⁺)VI | 0,37 | 0,52 | 312 | 6.8 |
| | a_2 - C5(Fe ³⁺)IV | 0,27 | 0,02 | 449 | 12,3 |
| | <i>c</i> ₂ - C6(Fe ³⁺)IV | 0,26 | 0,32 | 366 | 13,8 |
| | $e-C7(Fe^{3+})V$ | 0,29 | 2,21 | 400 | 8,9 |

 $\label{eq:table_$

 ${\bf Table} \ {\bf 2}-{\bf Magnetic} \ {\bf properties} \ {\bf of} \ {\bf anisotropic} \ {\bf hexagonal} \ {\bf ferrites}$

| Ferrites | The specific satu- ration magnetiza- tion σ_{s} , Am ² /kg | Specific residual magnetization σ_r , Am ² /kg | Coercive force <i>Hc</i> , kA/m | The Curie tem- perature, K |
|------------------------------------|--|--|------------------------------------|-------------------------------|
| BaFe ₁₂ O ₁₉ | 64,61 | 40,13 | 273 | 720 |
| BaFe10Al2O19 | 21,68 | 12,13 | 532,3 | 543 |

The area of the hysteresis loop in the ferrite $BaFe_{12}O_{19}$ reduced by 8 times, and M_r / M_s 1.5 times (from -0.641 to 0.560). Considering that the shape of the hysteresis loop depends on how the occur reversal processes, can be assumed that the substitution of Fe³⁺ ions Al³⁺ ions changing mechanical stresses or internal parameters of the spin-lattice relaxation. This is indicated by the changes the isomer shift and quadrupole splitting of the Mössbauer spectra for the sublattices aand c, which showed that the electronic structure transforms of Fe³⁺ ions as a result of changes in their local environment (Table 1). Indeed each cation located in the octahedral position 12k surrounded by six nearest neighbors (cations) from the trigonal axis of symmetry. If Al^{3+} ion occupies the tetrahedral position, the presence of the energy of the spin-lattice interaction is necessary to minimize its orderly arrangement of ions Al³⁺. In this case, Fe³⁺ ions are displaced position, which reinforce the relationship between the spin Fe³⁺ ions and the corresponding axis of the crystal field. It is the reason additional local the crystal lattice deformations and increase the anisotropy field in the ferrite BaFe10Al2O12 experimentally observed. Thus, the joint variation of local characteristics of the hexagonal ferrite with the M structure resulting from substitution of Fe³⁺ ions from the ion radius close Al³⁺ ions allows obtaining a material with combination of properties necessary for device millimeter waveband.

4. CONCLUSION

The investigations of hexagonal ferrites $BaFe_{12-x}Al_xO_{12}$ with isomorphic substitution of iron aluminum, allowed to establish a preferential distribution of ions Al^{3+} on the border of spinel and hexagonal blocks structure. Such localization of diamagnetic Al^{3+} contributed to a decrease in magnetization of hexagonal barium ferrite, a change in its magnetic structure and magnetic anisotropy field.

The results can be used in the selection of optimum compositions hexagonal ferrites substrate millimeter wave devices.

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